

UDC 512.5

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BRANCHING LAW FOR THE FINITE SUBGROUPS OF $\mathrm{SL}_4\mathbb{C}$ AND THE RELATED GENERALIZED POINCARÉ POLYNOMIALS

ЗАКОН ГАЛУЖЕННЯ ДЛЯ СКІНЧЕННИХ ПІДГРУП $\mathrm{SL}_4\mathbb{C}$ ТА ВІДПОВІДНІ УЗАГАЛЬНЕНІ ПОЛІНОМИ ПУАНКАРЕ

Within the framework of McKay correspondence we determine, for every finite subgroup Γ of $\mathrm{SL}_4\mathbb{C}$, how the finite-dimensional irreducible representations of $\mathrm{SL}_4\mathbb{C}$ decompose under the action of Γ .

Let \mathfrak{h} be a Cartan subalgebra of $\mathfrak{sl}_4\mathbb{C}$ and let $\varpi_1, \varpi_2, \varpi_3$ be the corresponding fundamental weights. For $(p, q, r) \in \mathbb{N}^3$, the restriction $\pi_{p,q,r}|_\Gamma$ of the irreducible representation $\pi_{p,q,r}$ of highest weight $p\varpi_1 + q\varpi_2 + r\varpi_3$ of $\mathrm{SL}_4\mathbb{C}$ decomposes as $\pi_{p,q,r}|_\Gamma = \bigoplus_{i=0}^l m_i(p, q, r)\gamma_i$, where $\{\gamma_0, \dots, \gamma_l\}$ is the set of equivalence classes of irreducible finite-dimensional complex representations of Γ . We determine the multiplicities $m_i(p, q, r)$ and prove that the series

$$P_\Gamma(t, u, w)_i = \sum_{p=0}^{\infty} \sum_{q=0}^{\infty} \sum_{r=0}^{\infty} m_i(p, q, r) t^p u^q w^r$$

are rational functions.

This generalizes the results of Kostant for $\mathrm{SL}_2\mathbb{C}$ and the results of our preceding works for $\mathrm{SL}_3\mathbb{C}$.

У рамках відповідності Маккея для кожної скінченної підгрупи Γ групи $\mathrm{SL}_4\mathbb{C}$ визначено, яким чином скінченно-вимірне незвіднє зображення $\mathrm{SL}_4\mathbb{C}$ розкладається під дією Γ .

Нехай \mathfrak{h} — картанова підалгебра $\mathfrak{sl}_4\mathbb{C}$, а $\varpi_1, \varpi_2, \varpi_3$ — відповідні фундаментальні ваги. Для $(p, q, r) \in \mathbb{N}^3$ зображення $\pi_{p,q,r}|_\Gamma$ незвідного зображення $\pi_{p,q,r}$ найбільшої ваги $p\varpi_1 + q\varpi_2 + r\varpi_3$ в $\mathrm{SL}_4\mathbb{C}$ розкладається у вигляді $\pi_{p,q,r}|_\Gamma = \bigoplus_{i=0}^l m_i(p, q, r)\gamma_i$, де $\{\gamma_0, \dots, \gamma_l\}$ — множина класів еквівалентності незвідних скінченновимірних комплексних зображень Γ . Визначено кратності $m_i(p, q, r)$ та доведено, що ряди

$$P_\Gamma(t, u, w)_i = \sum_{p=0}^{\infty} \sum_{q=0}^{\infty} \sum_{r=0}^{\infty} m_i(p, q, r) t^p u^q w^r$$

є раціональними функціями.

Це є узагальненням результатів Костанта для $\mathrm{SL}_2\mathbb{C}$, а також результатів наших попередніх робіт для $\mathrm{SL}_3\mathbb{C}$.

1. Introduction and results. Let Γ be a finite subgroup of $\mathrm{SL}_4\mathbb{C}$ and $\{\gamma_0, \dots, \gamma_l\}$ the set of equivalence classes of irreducible finite dimensional complex representations of Γ , where γ_0 is the trivial representation. The character associated to γ_j is denoted by χ_j .

Consider $\gamma: \Gamma \rightarrow \mathrm{SL}_4\mathbb{C}$ the natural 4-dimensional representation, and γ^* its contragredient representation. The character of γ is denoted by χ . By complete reducibility we get the decompositions

$$\forall j \in \llbracket 0, l \rrbracket: \gamma_j \otimes \gamma = \bigoplus_{i=0}^l a_{ij}^{(1)} \gamma_i, \quad \gamma_j \otimes (\gamma \wedge \gamma) = \bigoplus_{i=0}^l a_{ij}^{(2)} \gamma_i \quad \text{and} \quad \gamma_j \otimes \gamma^* = \bigoplus_{i=0}^l a_{ij}^{(3)} \gamma_i.$$

This defines the three following square matrices of $\mathbf{M}_{l+1}\mathbb{N}$:

$$A^{(1)} := \left(a_{ij}^{(1)} \right)_{(i,j) \in \llbracket 0, l \rrbracket^2}, \quad A^{(2)} := \left(a_{ij}^{(2)} \right)_{(i,j) \in \llbracket 0, l \rrbracket^2} \quad \text{and} \quad A^{(3)} := \left(a_{ij}^{(3)} \right)_{(i,j) \in \llbracket 0, l \rrbracket^2}.$$

Let \mathfrak{h} be a Cartan subalgebra of $\mathfrak{sl}_4\mathbb{C}$ and let $\varpi_1, \varpi_2, \varpi_3$ be the corresponding fundamental weights, and $V(p\varpi_1 + q\varpi_2 + r\varpi_3)$ the simple $\mathfrak{sl}_4\mathbb{C}$ -module of highest weight $p\varpi_1 + q\varpi_2 + r\varpi_3$

with $(p, q, r) \in \mathbb{N}^3$. Then we get an irreducible representation $\pi_{p,q,r}: \mathbf{SL}_4\mathbb{C} \rightarrow \mathbf{GL}(V(p\varpi_1 + q\varpi_2 + r\varpi_3))$. The restriction of $\pi_{p,q,r}$ to the subgroup Γ is a representation of Γ , and by complete reducibility, we get the decomposition

$$\pi_{p,q,r}|_\Gamma = \bigoplus_{i=0}^l m_i(p, q, r) \gamma_i,$$

where the $m_i(p, q, r)$'s are non negative integers. Let $\mathcal{E} := (e_0, \dots, e_l)$ be the canonical basis of \mathbb{C}^{l+1} , and

$$v_{p,q,r} := \sum_{i=0}^l m_i(p, q, r) e_i \in \mathbb{C}^{l+1}.$$

We have in particular $v_{0,0,0} = e_0$ as γ_0 is the trivial representation. Let us consider the vector

$$P_\Gamma(t, u, w) := \sum_{p=0}^{\infty} \sum_{q=0}^{\infty} \sum_{r=0}^{\infty} v_{p,q,r} t^p u^q w^r \in (\mathbb{C}[[t, u, w]])^{l+1},$$

and denote by $P_\Gamma(t, u, w)_j$ its j th coordinate in the basis \mathcal{E} , which is an element of $\mathbb{C}[[t, u, w]]$. Note that $P_\Gamma(t, u, w)$ can also be seen as a formal power series with coefficients in \mathbb{C}^{l+1} . The aim of this article is to prove the following theorem.

Theorem 1. *The coefficients of $P_\Gamma(t, u, w)$ are rational fractions in t, u, w , i.e., the formal power series $P_\Gamma(t, u, w)_i$ are rational functions*

$$P_\Gamma(t, u, w)_i = \frac{N_\Gamma(t, u, w)_i}{D_\Gamma(t, u, w)}, \quad i \in \llbracket 0, l \rrbracket,$$

where the $N_\Gamma(t, u, w)_i$'s and $D_\Gamma(t, u, w)$ are elements of $\mathbb{Q}[t, u, w]$.

The proof of this theorem uses a key-relation satisfied by $P_\Gamma(t, u, w)$ as well as a so-called inversion formula. Two essential ingredients are the decomposition of the tensor product of $\pi_{p,q,r}$ with the natural representation of $\mathbf{SL}_4\mathbb{C}$ and the simultaneous diagonalizability of certain matrices. The effective calculation of $P_\Gamma(t, u, w)$ then reduces to matrix multiplication.

In [2] we applied a similar method for $\mathbf{SL}_2\mathbb{C}$ — recovering thereby in a quite easy way the results obtained by Kostant in [6, 7], and by Gonzalez–Sprinberg and Verdier in [4] — and for $\mathbf{SL}_3\mathbb{C}$ in order to get explicit computations of the series for every finite subgroup of $\mathbf{SL}_3\mathbb{C}$.

The general framework of that study is the construction of a minimal resolution of singularities of the orbifold \mathbb{C}^n/Γ . It is related to the McKay correspondence (see [1, 3, 4]). For example, Gonzalez–Sprinberg and Verdier use in [4] a Poincaré series to construct explicitly minimal resolutions for singularities of $V = \mathbb{C}^2/\Gamma$ when Γ is a finite subgroup of $\mathbf{SL}_2\mathbb{C}$. To go further in this approach, our results for $\mathbf{SL}_4\mathbb{C}$ could be used to construct an explicit synthetic minimal resolution of singularities for orbifolds of the form \mathbb{C}^4/Γ where Γ is a finite subgroup of $\mathbf{SL}_4\mathbb{C}$.

2. Properties of the matrices $A^{(1)}, A^{(2)}, A^{(3)}$. In order to compute the series $P_\Gamma(t, u, w)$, we first establish here some properties of the matrices $A^{(1)}, A^{(2)}, A^{(3)}$. The first proposition essentially follows from the uniqueness of the decomposition of a representation as sum of irreducible representations.

Proposition 1. (i) $A^{(3)} = {}^t A^{(1)}$. (ii) $A^{(2)}$ is a symmetric matrix. (iii) $A^{(1)}$, $A^{(2)}$ and $A^{(3)}$ commute. In particular, $A^{(1)}$ is a normal matrix.

Proof. Since $a_{ij}^{(1)} = (\chi_i | \chi_{\gamma \otimes \gamma_j}) = \frac{1}{|\Gamma|} \sum_{g \in \Gamma} \overline{\chi_i(g)} \chi(g) \chi_j(g)$, we have $\gamma \otimes \gamma_j = \bigoplus_{i=0}^l a_{ij}^{(1)} \gamma_i$.

In the same way, $(\gamma \wedge \gamma) \otimes \gamma_j = \bigoplus_{i=0}^l a_{ij}^{(2)} \gamma_i$ and $\gamma^* \otimes \gamma_j = \bigoplus_{i=0}^l a_{ij}^{(3)} \gamma_i$.

Then

$$\begin{aligned} a_{ij}^{(3)} &= (\chi_i | \chi_{\gamma_j \otimes \gamma^*}) = \frac{1}{|\Gamma|} \sum_{g \in \Gamma} \overline{\chi_i(g)} \chi_j(g) \chi_{\gamma^*}(g) = \frac{1}{|\Gamma|} \sum_{g \in \Gamma} \overline{\chi_i(g)} \chi_j(g) \chi(g^{-1}) = \\ &= \frac{1}{|\Gamma|} \sum_{g \in \Gamma} \overline{\chi_i(g^{-1})} \chi_j(g^{-1}) \chi(g) = \frac{1}{|\Gamma|} \sum_{g \in \Gamma} \chi_i(g) \overline{\chi_j(g)} \chi(g) = a_{ji}^{(1)}, \end{aligned}$$

hence $A^{(3)} = {}^t A^{(1)}$.

We also have $(\gamma_j \otimes \gamma) \otimes \gamma^* = \left(\bigoplus_{i=0}^l a_{ij}^{(1)} \gamma_i \right) \otimes \gamma^* = \bigoplus_{i=0}^l a_{ij}^{(1)} \left(\bigoplus_{k=0}^l a_{ki}^{(3)} \gamma_k \right) = \bigoplus_{k=0}^l \left(\sum_{i=0}^l a_{ki}^{(3)} a_{ij}^{(1)} \right) \gamma_k$ and $\gamma \otimes (\gamma_j \otimes \gamma^*) = \gamma \otimes \left(\bigoplus_{i=0}^l a_{ij}^{(3)} \gamma_i \right) = \bigoplus_{i=0}^l a_{ij}^{(3)} \left(\bigoplus_{k=0}^l a_{ki}^{(1)} \gamma_k \right) = \bigoplus_{k=0}^l \left(\sum_{i=0}^l a_{ki}^{(1)} a_{ij}^{(3)} \right) \gamma_k$, hence $A^{(3)} A^{(1)} = A^{(1)} A^{(3)}$. The proofs of the other statements are the same.

Since $A^{(1)}$, $A^{(2)}$, $A^{(3)}$ are normal, we know that they are diagonalizable with eigenvectors forming an orthogonal basis. Now we will diagonalize these matrices by using the character table of the group Γ . Let us denote by $\{C_0, \dots, C_l\}$ the set of conjugacy classes of Γ , and for any $j \in \llbracket 0, l \rrbracket$, let g_j be an element of C_j . So the character table of Γ is the matrix $T_\Gamma \in \mathbf{M}_{l+1}\mathbb{C}$ defined by $(T_\Gamma)_{i,j} := \chi_i(g_j)$.

Proposition 2. (i) For $k \in \llbracket 0, l \rrbracket$, set $w_k := (\chi_0(g_k), \dots, \chi_l(g_k)) \in \mathbb{C}^{l+1}$. Then w_k is an eigenvector of $A^{(3)}$ associated to the eigenvalue $\chi(g_k)$. Similarly, w_k is an eigenvector of $A^{(1)}$ associated to the eigenvalue $\overline{\chi(g_k)}$.

(ii) For $k \in \llbracket 0, l \rrbracket$, w_k is an eigenvector of $A^{(2)}$ associated to the eigenvalue $\frac{1}{2}(\chi(g_k)^2 + \chi(g_k^2))$.

Proof. From the relation $\gamma_i \otimes \gamma = \sum_{j=0}^l a_{ji}^{(1)} \gamma_j$, we get $\chi_i \chi = \chi_{\gamma_i \otimes \gamma} = \sum_{j=0}^l a_{ji}^{(1)} \chi_j$. By evaluating this on g_k , we obtain $\chi_i(g_k) \chi(g_k) = \sum_{j=0}^l a_{ji}^{(1)} \chi_j(g_k) = \sum_{j=0}^l a_{ij}^{(3)} \chi_j(g_k)$ according to Proposition 1. So w_k is an eigenvector of $A^{(3)}$ associated to the eigenvalue $\chi(g_k)$. The method is similar for the other results.

As the w_j 's are the columns of T_Γ , which are always orthogonal, the matrix T_Γ is invertible and the family $\mathcal{W} := (w_0, \dots, w_l)$ is a common basis of eigenvectors of $A^{(1)}$, $A^{(2)}$ and $A^{(3)}$. Then $\Lambda^{(1)} := T_\Gamma^{-1} A^{(1)} T_\Gamma$, $\Lambda^{(2)} := T_\Gamma^{-1} A^{(2)} T_\Gamma$ and $\Lambda^{(3)} := T_\Gamma^{-1} A^{(3)} T_\Gamma$ are diagonal matrices, with $\Lambda_{jj}^{(1)} = \overline{\chi(g_j)}$, $\Lambda_{jj}^{(2)} = \frac{1}{2}(\chi(g_j)^2 - \chi(g_j^2))$ and $\Lambda_{jj}^{(3)} = \chi(g_j)$.

Now, we make use of the Clebsch–Gordan formula

$$\pi_{1,0,0} \otimes \pi_{p,q,r} = \pi_{p+1,q,r} \oplus \pi_{p,q,r-1} \oplus \pi_{p-1,q+1,r} \oplus \pi_{p,q-1,r+1},$$

$$\pi_{0,1,0} \otimes \pi_{p,q,r} = \pi_{p,q+1,r} \oplus \pi_{p,q-1,r} \oplus \pi_{p+1,q-1,r+1} \oplus \pi_{p-1,q+1,r-1} \oplus \pi_{p-1,q,r+1} \oplus \pi_{p+1,q,r-1}, \quad (1)$$

$$\pi_{0,0,1} \otimes \pi_{p,q,r} = \pi_{p,q,r+1} \oplus \pi_{p-1,q,r} \oplus \pi_{p,q+1,r-1} \oplus \pi_{p+1,q-1,r}.$$

Proposition 3. *The vectors $v_{p,q,r}$ satisfy the following recurrence relations:*

$$A^{(1)}v_{p,q,r} = v_{p+1,q,r} + v_{p,q,r-1} + v_{p-1,q+1,r} + v_{p,q-1,r+1},$$

$$A^{(2)}v_{p,q,r} = v_{p,q+1,r} + v_{p,q-1,r} + v_{p+1,q-1,r+1} + v_{p-1,q+1,r-1} + v_{p-1,q,r+1} + v_{p+1,q,r-1},$$

$$A^{(3)}v_{p,q,r} = v_{p,q,r+1} + v_{p-1,q,r} + v_{p,q+1,r-1} + v_{p+1,q-1,r}.$$

Proof. The definition of $v_{p,q,r}$ reads $v_{p,q,r} = \sum_{i=0}^l m_i(p, q, r) e_i$, thus $A^{(1)}v_{p,q,r} = \sum_{i=0}^l \left(\sum_{j=0}^l m_j(p, q, r) a_{ij}^{(1)} \right) e_i$. Now

$$(\pi_{1,0,0} \otimes \pi_{p,q,r})|_\Gamma = \pi_{p,q,r}|_\Gamma \otimes \gamma = \sum_{j=0}^l m_j(p, q, r) \gamma_j \otimes \gamma = \sum_{i=0}^l \left(\sum_{j=0}^l m_j(p, q, r) a_{ij}^{(1)} \right) \gamma_i$$

and

$$\begin{aligned} & \pi_{p+1,q,r}|_\Gamma + \pi_{p,q,r-1}|_\Gamma + \pi_{p-1,q+1,r}|_\Gamma + \pi_{p,q-1,r+1}|_\Gamma = \\ & = \sum_{i=0}^l (m_i(p+1, q, r) + m_i(p, q, r-1) + m_i(p-1, q+1, r) + m_i(p, q-1, r+1)) \gamma_i. \end{aligned}$$

By uniqueness,

$$\begin{aligned} & \sum_{j=0}^l m_j(p, q, r) a_{ij}^{(1)} = m_i(p+1, q, r) + m_i(p, q, r-1) + \\ & + m_i(p-1, q+1, r) + m_i(p, q-1, r+1). \end{aligned}$$

3. The series $P_\Gamma(t, u, w)$ is a rational function. This section is mainly devoted to the proof of Theorem 1.

3.1. A key-relation satisfied by the series $P_\Gamma(t, u, w)$.

Proposition 4. Set

$$\begin{aligned} J(t, u, w) := & (1 - u^2)((1 + ut^2)(1 + uw^2) - tw(1 + u^2))I_n + twu(1 - u^2)A^{(2)} - \\ & - tu(1 + uw^2)(A^{(3)} - uA^{(1)}) - wu(1 + ut^2)(A^{(1)} - uA^{(3)}). \end{aligned}$$

Then the series $P_\Gamma(t, u, w)$ satisfies the following relation:

$$\begin{aligned} & J(t, u, w) v_{0,0,0} = \\ & = (1 - tA^{(1)} + t^2 A^{(2)} - t^3 A^{(3)} + t^4) (1 - wA^{(3)} + w^2 A^{(2)} - w^3 A^{(1)} + w^4) \times \\ & \times ((1 + u^2)(1 - u^2)^2 - u(1 - u^2)^2 A^{(2)} + u^2(A^{(1)} - uA^{(3)})(A^{(3)} - uA^{(1)})) P_\Gamma(t, u, w). \end{aligned}$$

Proof. Set $x := P_{\Gamma}(t, u, w)$. Set also $v_{p,q,-1} := 0$, $v_{p,-1,r} := 0$ and $v_{-1,q,r} := 0$ for $(p, q, r) \in \mathbb{N}^3$, such that, according to the Clebsch–Gordan formula, the formulae of the preceding corollary are still true for $(p, q, r) \in \mathbb{N}^3$. So we have (by denoting $\sum_{p=0}^{\infty} \sum_{q=0}^{\infty} \sum_{r=0}^{\infty}$ by \sum_{pqr})

$$\begin{aligned} & (1 - wA^{(3)} + w^2A^{(2)} - w^3A^{(1)} + w^4)x = \\ & = \sum_{pqr} v_{p,q,r} t^p u^q w^r - \sum_{pqr} (v_{p,q,r+1} + v_{p-1,q,r} + v_{p,q+1,r-1} + v_{p+1,q-1,r}) t^p u^q w^{r+1} + \\ & + \sum_{pqr} (v_{p,q+1,r} + v_{p,q-1,r} + v_{p+1,q-1,r+1} + v_{p-1,q+1,r-1} + v_{p-1,q,r+1} + v_{p+1,q,r-1}) t^p u^q w^{r+2} - \\ & - \sum_{pqr} (v_{p+1,q,r} + v_{p,q,r-1} + v_{p-1,q+1,r} + v_{p,q-1,r+1}) t^p u^q w^{r+3} + \sum_{pqr} v_{p,q,r} t^p u^q w^{r+4}, \end{aligned}$$

hence

$$\begin{aligned} & (1 - wA^{(3)} + w^2A^{(2)} - w^3A^{(1)} + w^4)x = \\ & = (1 - tw + uw^2 - t^{-1}uw) \sum_{p=0}^{\infty} \sum_{q=0}^{\infty} v_{p,q,0} t^p u^q + t^{-1}uw \sum_{q=0}^{\infty} v_{0,q,0} u^q. \end{aligned} \quad (2)$$

In the same way (by denoting $\sum_{p=0}^{\infty} \sum_{q=0}^{\infty}$ by \sum_{pq})

$$\begin{aligned} & (1 - tA^{(1)} + t^2A^{(2)} - t^3A^{(3)} + t^4) \sum_{p=0}^{\infty} \sum_{q=0}^{\infty} v_{p,q,0} t^p u^q = \\ & = \sum_{pq} v_{p,q,0} t^p u^q - \sum_{pq} (v_{p+1,q,0} + v_{p-1,q+1,0} + v_{p,q-1,1}) t^{p+1} u^q + \\ & + \sum_{pq} (v_{p,q+1,0} + v_{p,q-1,0} + v_{p+1,q-1,1} + v_{p-1,q,1}) t^{p+2} u^q - \\ & - \sum_{pq} (v_{p,q,1} + v_{p-1,q,0} + v_{p+1,q-1,0}) t^{p+3} u^q + \sum_{pq} v_{p,q,0} t^{p+4} u^q, \end{aligned}$$

hence

$$\begin{aligned} & (1 - tA^{(1)} + t^2A^{(2)} - t^3A^{(3)} + t^4) \sum_{p=0}^{\infty} \sum_{q=0}^{\infty} v_{p,q,0} t^p u^q = \\ & = (1 + t^2u) \sum_{q=0}^{\infty} v_{0,q,0} u^q - tu \sum_{q=0}^{\infty} v_{0,q,1} u^q. \end{aligned} \quad (3)$$

Moreover, we have

$$\begin{aligned}
& (1 - tA^{(1)} + t^2A^{(2)} - t^3A^{(3)} + t^4) \sum_{q=0}^{\infty} v_{0,q,0}u^q = \\
& = \sum_{q=0}^{\infty} v_{0,q,0}u^q - \sum_{q=0}^{\infty} (v_{1,q,0} + v_{0,q-1,1})tu^q + \sum_{q=0}^{\infty} (v_{0,q+1,0} + v_{0,q-1,0} + v_{1,q-1,1})t^2u^q - \\
& \quad - \sum_{q=0}^{\infty} (v_{0,q,1} + v_{1,q-1,0})t^3u^q + \sum_{q=0}^{\infty} v_{0,q,0}t^4u^q,
\end{aligned}$$

hence

$$\begin{aligned}
& (1 - tA^{(1)} + t^2A^{(2)} - t^3A^{(3)} + t^4) \sum_{p=0}^{\infty} v_{0,p,0}u^p = \\
& = (1 + t^4 + t^2u^{-1} + t^2u) \sum_{q=0}^{\infty} v_{0,q,0}u^q - t^2u^{-1}v_{0,0,0} - (t + t^3u) \sum_{q=0}^{\infty} v_{1,q,0}u^q - \\
& \quad - (tu + t^3) \sum_{q=0}^{\infty} v_{0,q,1}u^q + t^2u \sum_{q=0}^{\infty} v_{1,q,1}u^q. \tag{4}
\end{aligned}$$

By combining equations (2), (3) and (4), we get

$$\begin{aligned}
& (1 - tA^{(1)} + t^2A^{(2)} - t^3A^{(3)} + t^4)(1 - wA^{(3)} + w^2A^{(2)} - w^3A^{(3)} + w^4)x = \\
& = (1 - tA^{(1)} + t^2A^{(2)} - t^3A^{(3)} + t^4) \times \\
& \quad \times \left((1 - tw + uw^2 - t^{-1}uw) \sum_{pq} v_{p,q,0}t^p u^q + t^{-1}uw \sum_{q=0}^{\infty} v_{0,q,0}u^q \right) = \\
& = (1 - tw + uw^2 - t^{-1}uw) \left((1 + t^2u) \sum_{q=0}^{\infty} v_{0,q,0}u^q - tu \sum_{q=0}^{\infty} v_{0,q,1}u^q \right) + \\
& \quad + (1 + t^4 + t^2u^{-1} + t^2u)t^{-1}uw \sum_{q=0}^{\infty} v_{0,q,0}u^q - twv_{0,0,0} - (1 + t^2u)uw \sum_{q=0}^{\infty} v_{1,q,0}u^q - \\
& \quad - (u + t^2)uw \sum_{q=0}^{\infty} v_{0,q,1}u^q + tu^2w \sum_{q=0}^{\infty} v_{1,q,1}u^q,
\end{aligned}$$

hence

$$(1 - tA^{(1)} + t^2A^{(2)} - t^3A^{(3)} + t^4)(1 - wA^{(3)} + w^2A^{(2)} - w^3A^{(1)} + w^4)x =$$

$$\begin{aligned}
 &= (1 + ut^2)(1 + uw^2) \sum_{q=0}^{\infty} v_{0,q,0} u^q - tu(1 + uw^2) \sum_{q=0}^{\infty} v_{0,q,1} u^q - \\
 &\quad - wu(1 + ut^2) \sum_{q=0}^{\infty} v_{1,q,0} u^q - twv_{0,0,0} + tu^2 w \sum_{q=0}^{\infty} v_{1,q,1} u^q. \tag{5}
 \end{aligned}$$

Besides, we have the two following equations:

$$A^{(1)} \sum_{q=0}^{\infty} v_{0,q,0} u^q = \sum_{q=0}^{\infty} v_{1,q,0} u^q + u \sum_{q=0}^{\infty} v_{0,q,1} u^q, \tag{6}$$

and

$$A^{(3)} \sum_{q=0}^{\infty} v_{0,q,0} u^q = \sum_{q=0}^{\infty} v_{0,q,1} u^q + u \sum_{q=0}^{\infty} v_{1,q,0} u^q. \tag{7}$$

From these two equations, we deduce

$$\sum_{q=0}^{\infty} v_{0,q,1} u^q = (1 - u^2)^{-1} (A^{(3)} - uA^{(1)}) \sum_{q=0}^{\infty} v_{0,q,0} u^q. \tag{8}$$

Now, we have

$$A^{(1)} \sum_{q=0}^{\infty} v_{0,q,1} u^q = \sum_{q=0}^{\infty} v_{1,q,1} u^q + \sum_{q=0}^{\infty} v_{0,q,0} u^q + u \sum_{q=0}^{\infty} v_{0,q,2} u^q, \tag{9}$$

and

$$A^{(3)} \sum_{q=0}^{\infty} v_{0,q,1} u^q = \sum_{q=0}^{\infty} v_{0,q,2} u^q + u^{-1} \sum_{q=0}^{\infty} v_{0,q,0} u^q + u \sum_{q=0}^{\infty} v_{1,q,1} u^q - u^{-1} v_{0,0,0}, \tag{10}$$

therefore

$$\sum_{q=0}^{\infty} v_{1,q,1} u^q = (1 - u^2)^{-1} (A^{(1)} - uA^{(3)}) \sum_{q=0}^{\infty} v_{0,q,1} u^q - (1 - u^2)^{-1} v_{0,0,0}.$$

So, according to equation (8), we deduce

$$\sum_{q=0}^{\infty} v_{1,q,1} u^q = (1 - u^2)^{-2} (A^{(1)} - uA^{(3)}) (A^{(3)} - uA^{(1)}) \sum_{q=0}^{\infty} v_{0,q,0} u^q - (1 - u^2)^{-1} v_{0,0,0}. \tag{11}$$

By using equation (11), we may write equation (5) as

$$\begin{aligned}
 &(1 - tA^{(1)} + t^2 A^{(2)} - t^3 A^{(3)} + t^4)(1 - wA^{(3)} + w^2 A^{(2)} - w^3 A^{(1)} + w^4)x = \\
 &= \left((1 + ut^2)(1 + uw^2) + tu^2 w(1 - u^2)^{-2} (A^{(1)} - uA^{(3)}) (A^{(3)} - uA^{(1)}) \right) \sum_{q=0}^{\infty} v_{0,q,0} u^q -
 \end{aligned}$$

$$-tu(1+uw^2)\sum_{q=0}^{\infty}v_{0,q,1}u^q-wu(1+ut^2)\sum_{q=0}^{\infty}v_{1,q,0}u^q-(tw+tu^2w(1-u^2)^{-1})v_{0,0,0}. \quad (12)$$

From equations (6) and (7), we also deduce

$$\sum_{q=0}^{\infty}v_{1,q,0}u^q=(1-u^2)^{-1}(A^{(1)}-uA^{(3)})\sum_{q=0}^{\infty}v_{0,q,0}u^q. \quad (13)$$

So, by using equations (8) and (13), we obtain

$$\begin{aligned} &(1-tA^{(1)}+t^2A^{(2)}-t^3A^{(3)}+t^4)(1-wA^{(3)}+w^2A^{(2)}-w^3A^{(1)}+w^4)x= \\ &= \left((1+ut^2)(1+uw^2)-tu(1+uw^2)(1-u^2)^{-1}(A^{(3)}-uA^{(1)})- \right. \\ &\quad \left. -wu(1+ut^2)(1-u^2)^{-1}(A^{(1)}-uA^{(3)})+ \right. \\ &\quad \left. +tu^2w(1-u^2)^{-2}(A^{(1)}-uA^{(3)})(A^{(3)}-uA^{(1)}) \right) \sum_{q=0}^{\infty}v_{0,q,0}u^q- \\ &\quad -(tw+tu^2w(1-u^2)^{-1})v_{0,0,0}, \end{aligned} \quad (14)$$

i.e., by multiplying (14) by $(1-u^2)^2$,

$$\begin{aligned} &(1-u^2)^2(1-tA^{(1)}+t^2A^{(2)}-t^3A^{(3)}+t^4)(1-wA^{(3)}+w^2A^{(2)}-w^3A^{(1)}+w^4)x= \\ &= \left((1-u^2)^2(1+ut^2)(1+uw^2)-tu(1+uw^2)(1-u^2)(A^{(3)}-uA^{(1)})- \right. \\ &\quad \left. -wu(1+ut^2)(1-u^2)(A^{(1)}-uA^{(3)})+tu^2w(A^{(1)}-uA^{(3)})(A^{(3)}-uA^{(1)}) \right) \sum_{q=0}^{\infty}v_{0,q,0}u^q- \\ &\quad -(tw(1-u^2)^2+tu^2w(1-u^2))v_{0,0,0}. \end{aligned} \quad (15)$$

Consider now the following equation:

$$A^{(2)}\sum_{q=0}^{\infty}v_{0,q,0}u^q=u^{-1}\sum_{q=0}^{\infty}v_{0,q,0}u^q+u\sum_{q=0}^{\infty}v_{0,q,0}u^q+u\sum_{q=0}^{\infty}v_{1,q,1}u^q-u^{-1}v_{0,0,0}. \quad (16)$$

Then, according to equation (11), we have

$$\begin{aligned} &A^{(2)}\sum_{q=0}^{\infty}v_{0,q,0}u^q=u^{-1}\sum_{q=0}^{\infty}v_{0,q,0}u^q+u\sum_{q=0}^{\infty}v_{0,q,0}u^q+ \\ &+u(1-u^2)^{-2}(A^{(1)}-uA^{(3)})(A^{(3)}-uA^{(1)})\sum_{q=0}^{\infty}v_{0,q,0}u^q-u(1-u^2)^{-1}v_{0,0,0}-u^{-1}v_{0,0,0}, \end{aligned}$$

i.e.,

$$\begin{aligned} & \left(A^{(2)} - u^{-1} - u - u(1-u^2)^{-2}(A^{(1)} - uA^{(3)})(A^{(3)} - uA^{(1)}) \right) \sum_{q=0}^{\infty} v_{0,q,0} u^q = \\ & = -(u(1-u^2)^{-1} + u^{-1})v_{0,0,0}. \end{aligned} \quad (17)$$

This last equation reads

$$\begin{aligned} & \left(-u(1-u^2)^2 A^{(2)} + (1+u^2)(1-u^2)^2 + u^2(A^{(1)} - uA^{(3)})(A^{(3)} - uA^{(1)}) \right) \sum_{q=0}^{\infty} v_{0,q,0} u^q = \\ & = (1-u^2)v_{0,0,0}. \end{aligned} \quad (18)$$

Now, by using equations (15) and (18), we get

$$\begin{aligned} & (1-u^2)^2(1-tA^{(1)} + t^2A^{(2)} - t^3A^{(3)} + t^4)(1-wA^{(3)} + w^2A^{(2)} - w^3A^{(1)} + w^4) \times \\ & \times \left(-u(1-u^2)^2 A^{(2)} + (1+u^2)(1-u^2)^2 + u^2(A^{(1)} - uA^{(3)})(A^{(3)} - uA^{(1)}) \right) x = \\ & = -tw(1-u^2) \left(-u(1-u^2)^2 A^{(2)} + (1+u^2)(1-u^2)^2 + \right. \\ & \left. + u^2(A^{(1)} - uA^{(3)})(A^{(3)} - uA^{(1)}) \right) v_{0,0,0} + \\ & + \left((1-u^2)^2(1+ut^2)(1+uw^2) - tu(1+uw^2)(1-u^2)(A^{(3)} - uA^{(1)}) - \right. \\ & \left. - wu(1+ut^2)(1-u^2)(A^{(1)} - uA^{(3)}) + tu^2w(A^{(1)} - uA^{(3)})(A^{(3)} - uA^{(1)}) \right) (1-u^2)v_{0,0,0}, \quad (19) \end{aligned}$$

i.e., after simplification by $(1-u^2)$,

$$\begin{aligned} & (1-tA^{(1)} + t^2A^{(2)} - t^3A^{(3)} + t^4)(1-wA^{(3)} + w^2A^{(2)} - w^3A^{(1)} + w^4) \times \\ & \times \left((1+u^2)(1-u^2)^2 - u(1-u^2)^2 A^{(2)} + u^2(A^{(1)} - uA^{(3)})(A^{(3)} - uA^{(1)}) \right) x = \\ & = \left((1-u^2)((1+ut^2)(1+uw^2) - tw(1+u^2)) + twu(1-u^2)A^{(2)} - \right. \\ & \left. - tu(1+uw^2)(A^{(3)} - uA^{(1)}) - wu(1+ut^2)(A^{(1)} - uA^{(3)}) \right) v_{0,0,0}. \quad (20) \end{aligned}$$

Proposition 4 is proved.

3.2. An inversion formula. In order to inverse the relation obtained in Proposition 4 and get an explicit expression for $P_T(t, u, w)$, we need the rational function f defined by

$$f : \mathbb{C}^3 \rightarrow \mathbb{C}(t, u, w),$$

$$(d_1, d_2, d_3) \mapsto (1 - td_1 + t^2 d_2 - t^3 d_3 + t^4)^{-1} (1 - wd_3 + w^2 d_2 - w^3 d_1 + w^4)^{-1} \times \\ \times \left((1 + u^2)(1 - u^2)^2 - u(1 - u^2)^2 d_2 + u^2(d_1 - ud_3)(d_3 - ud_1) \right)^{-1}.$$

According to Proposition 4, we may write

$$J(t, u, w) v_{0,0,0} = \\ = T_\Gamma \left(1 - t\Lambda^{(1)} + t^2\Lambda^{(2)} - t^3\Lambda^{(3)} + t^4 \right) \left(1 - w\Lambda^{(3)} + w^2\Lambda^{(2)} - w^3\Lambda^{(1)} + w^4 \right) \times \\ \times \left((1 + u^2)(1 - u^2)^2 - u(1 - u^2)^2 \Lambda^{(2)} + u^2(\Lambda^{(1)} - u\Lambda^{(3)})(\Lambda^{(3)} - u\Lambda^{(1)}) \right) T_\Gamma^{-1} P_\Gamma(t, u, w).$$

We deduce that

$$P_\Gamma(t, u, w) = T_\Gamma \Delta(t, u, w) T_\Gamma^{-1} J(t, u, w) v_{0,0,0} = \\ = (T_\Gamma \Delta(t, u, w) T_\Gamma) (T_\Gamma^{-2} J(t, u, w) v_{0,0,0}), \quad (21)$$

where $\Delta(t, u, w) \in \mathbf{M}_{l+1}\mathbb{C}(t, u, w)$ is the diagonal matrix defined by

$$\Delta(t, u, w)_{jj} = f(\Lambda_{jj}^{(1)}, \Lambda_{jj}^{(2)}, \Lambda_{jj}^{(3)}) = f \left(\overline{\chi(g_j)}, \frac{1}{2}(\chi(g_j)^2 - \chi(g_j^2)), \chi(g_j) \right).$$

This last formula proves Theorem 1.

Remark 1. The Poincaré series $\hat{P}_\Gamma(t)$ of the algebra of invariants $\mathbb{C}[z_1, z_2, z_3, z_4]^\Gamma$ is given by

$$\hat{P}_\Gamma(t) = P_\Gamma(t, 0, 0)_0 = P_\Gamma(0, 0, t)_0.$$

3.3. Remark for $\mathbf{SL}_n\mathbb{C}$. In this section, we consider an integer $n \geq 2$ and a subgroup Γ of $\mathbf{SL}_n\mathbb{C}$. As in Section 1, let $\{\gamma_0, \dots, \gamma_l\}$ be the set of equivalence classes of irreducible finite dimensional complex representations of Γ , where γ_0 is the trivial representation. The character associated to γ_j is denoted by χ_j .

Consider $\gamma: \Gamma \rightarrow \mathbf{SL}_n\mathbb{C}$ the natural n -dimensional representation, and χ its character. By complete reducibility we get the decomposition $\gamma_j \otimes \gamma = \bigoplus_{i=0}^l a_{ij}^{(1)} \gamma_i$ for every $j \in \llbracket 0, l \rrbracket$, and we set $A^{(1)} := (a_{ij}^{(1)})_{(i,j) \in \llbracket 0, l \rrbracket^2} \in \mathbf{M}_{l+1}\mathbb{N}$.

Let \mathfrak{h} be a Cartan subalgebra of $\mathfrak{sl}_n\mathbb{C}$ and let $\varpi_1, \dots, \varpi_{n-1}$ be the corresponding fundamental weights, and $V(p_1\varpi_1 + \dots + p_{n-1}\varpi_{n-1})$ the simple $\mathfrak{sl}_n\mathbb{C}$ -module of highest weight $p_1\varpi_1 + \dots + p_{n-1}\varpi_{n-1}$ with $\mathbf{p} := (p_1, \dots, p_{n-1}) \in \mathbb{N}^{n-1}$. Then we get an irreducible representation $\pi_{\mathbf{p}}: \mathbf{SL}_n\mathbb{C} \rightarrow \mathbf{GL}(V(p_1\varpi_1 + \dots + p_{n-1}\varpi_{n-1}))$. The restriction of $\pi_{\mathbf{p}}$ to the subgroup Γ is a representation of Γ , and by complete reducibility, we get the decomposition $\pi_{\mathbf{p}}|_\Gamma = \bigoplus_{i=0}^l m_i(\mathbf{p}) \gamma_i$, where the $m_i(\mathbf{p})$'s are non negative integers. Let $\mathcal{E} := (e_0, \dots, e_l)$ be the canonical basis of \mathbb{C}^{l+1} , and

$$v_{\mathbf{p}} := \sum_{i=0}^l m_i(\mathbf{p}) e_i \in \mathbb{C}^{l+1}.$$

As γ_0 is the trivial representation, we have $v_0 = e_0$. Let us consider the vector (with elements of

$\mathbb{C}[[t_1, \dots, t_{n-1}]] = \mathbb{C}[[t]]$ as coefficients)

$$P_\Gamma(\mathbf{t}) := \sum_{\mathbf{p} \in \mathbb{N}^{n-1}} v_{\mathbf{p}} \mathbf{t}^{\mathbf{p}} \in (\mathbb{C}[[t]])^{l+1},$$

and denote by $P_\Gamma(\mathbf{t})_j$ its j th coordinate in the basis \mathcal{E} .

Given the results from Kostant [6, 7] for $\mathbf{SL}_2\mathbb{C}$ and our results [2] about $\mathbf{SL}_3\mathbb{C}$, we then formulate the following conjecture:

Conjecture 1. *The coefficients of the vector $P_\Gamma(\mathbf{t})$ are rational fractions in \mathbf{t} , i.e., the formal power series $P_\Gamma(\mathbf{t})_i$ are rational functions*

$$P_\Gamma(\mathbf{t})_i := \frac{N_\Gamma(\mathbf{t})_i}{D_\Gamma(\mathbf{t})}, \quad i \in \llbracket 0, l \rrbracket,$$

where the $N_\Gamma(\mathbf{t})_i$'s and $D_\Gamma(\mathbf{t})$ are elements of $\mathbb{Q}[[t]]$.

4. An example of explicit computation. The classification of finite subgroups of $\mathbf{SL}_4\mathbb{C}$ is given in [5]. It consists in infinite series and 30 exceptional groups (types I, II, \dots, XXX). We give here an explicit computation of $P_\Gamma(t, u, w)$ for one of these exceptional groups. Consider the matrices

$$F_1 = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & j & 0 \\ 0 & 0 & 0 & j^2 \end{pmatrix}, \quad F'_2 = \frac{1}{3} \begin{pmatrix} 3 & 0 & 0 & 0 \\ 0 & -1 & 2 & 2 \\ 0 & 2 & -1 & 2 \\ 0 & 2 & 2 & -1 \end{pmatrix},$$

$$F'_3 = \frac{1}{4} \begin{pmatrix} -1 & \sqrt{15} & 0 & 0 \\ \sqrt{15} & 1 & 0 & 0 \\ 0 & 0 & 0 & 4 \\ 0 & 0 & 4 & 0 \end{pmatrix},$$

and the subgroup $\Gamma = \langle F_1, F'_2, F'_3 \rangle$ of $\mathbf{SL}_4\mathbb{C}$ (type II in [5]).

Here $l = 4$,

$$A^{(1)} = A^{(3)} = \begin{pmatrix} 0 & 0 & 0 & 1 & 0 \\ 0 & 0 & 1 & 1 & 1 \\ 0 & 1 & 0 & 1 & 1 \\ 1 & 1 & 1 & 1 & 1 \\ 0 & 1 & 1 & 1 & 2 \end{pmatrix} \quad \text{and} \quad A^{(2)} = \begin{pmatrix} 0 & 1 & 1 & 0 & 0 \\ 1 & 1 & 0 & 1 & 2 \\ 1 & 0 & 1 & 1 & 2 \\ 0 & 1 & 1 & 2 & 2 \\ 0 & 2 & 2 & 2 & 2 \end{pmatrix},$$

$\text{rank}(A^{(1)}) = \text{rank}(A^{(2)}) = 4$, and the eigenvalues of $A^{(1)}, A^{(2)}, A^{(3)}$ are $\Theta^{(1)} = \overline{\Theta^{(3)}} = (4, 0, -1, 1, -1), \Theta^{(2)} = (6, -2, 1, 0, 1)$.

According to formula (21), we get

$$D_\Gamma(t, u, w) = (w-1)^4 (u+1)^3 (u-1)^5 (t-1)^4 (w^2 + w + 1) (w^4 + w^3 + w^2 + w + 1) \times$$

$$\begin{aligned} & \times (w+1)^2 (u^4 + u^3 + u^2 + u + 1) (u^2 + u + 1)^2 (t^2 + t + 1) (t^4 + t^3 + t^2 + t + 1) (t + 1)^2 = \\ & = (u - 1)(u + 1)(u^2 + u + 1)\tilde{D}_\Gamma(t)\tilde{D}_\Gamma(u)\tilde{D}_\Gamma(w), \end{aligned}$$

with $\tilde{D}_\Gamma(t) = (t - 1)^4(t + 1)^2(t^2 + t + 1)(t^4 + t^3 + t^2 + t + 1)$. Moreover,

$$\widehat{P}_\Gamma(t) = \frac{t^8 - t^6 + t^4 - t^2 + 1}{t^{12} - 2t^{10} - t^9 + t^8 + t^7 + t^5 + t^4 - t^3 - 2t^2 + 1}.$$

Because of the too big size of the numerators $N_\Gamma(t, u, w)_i$'s, only the denominator is given in the text: all the numerators may be found on the web (<http://math.univ-lyon1.fr/homes-www/butin/>).

1. Bridgeland T., King A., Reid M. The McKay correspondence as an equivalence of derived categories // J. Amer. Math. Soc. – 2001. – **14**. – P. 535–554.
2. Butin F., Perets G. S. McKay correspondence and the branching law for finite subgroups of $\mathbf{SL}_3\mathbb{C}$ // J. Group Theory. – 2014. – **17**, Issue 2. – P. 191–251.
3. Gomi Y., Nakamura I., Shinoda K. Coinvariant algebras of finite subgroups of $\mathbf{SL}_3\mathbb{C}$ // Can. J. Math. – 2004. – **56**. – P. 495–528.
4. Gonzalez-Sprinberg G., Verdier J.-L. Construction géométrique de la correspondance de McKay // An. sci. de l’E. N. S., 4ème sér. – 1983. – **16**, № 3. – P. 409–449.
5. Hanany A., He Y.-H. A monograph on the classification of the discrete subgroups of $\mathbf{SU}(4)$ // JHEP. – 2001. – **27**.
6. Kostant B. The McKay correspondence, the Coxeter element and representation theory // SMF. Astérisque, hors sér. – 1985. – P. 209–255.
7. Kostant B. The Coxeter element and the branching law for the finite subgroups of $\mathbf{SU}(2)$ // Coxeter Legacy. – Providence, RI: Amer. Math. Soc., 2006. – P. 63–70.

Received 14.01.14,
after revision – 29.03.15